

# Estimating Molecular Electronic Chemical Potential and Hardness from Fragments' Addition Schemes

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A new scheme for the estimation of aggregate chemical potential and hardness is introduced and compared with the results of ab initio calculations for the aggregates as well as with the results obtained by employing various other combination schemes. Numerical results show that this new scheme provides better estimates of electronic chemical potential as well as comparable results for hardness with other addition schemes.

## 1. Introduction

Within the purview of the density functional theory,<sup>1–3</sup> the electronegativity  $\chi$  of an atom has been identified<sup>4</sup> with the negative of the chemical potential  $\mu$  of the electronic distribution, and the chemical hardness  $\eta$  has been defined as the partial derivative of the chemical potential with respect to the total number of electrons at constant external potential.<sup>5</sup> The establishment of a rigorous foundation for the popular concepts of electronegativity and hardness allows for the introduction of various combination schemes relating the global properties of the constituent fragments to the global properties of the resultant aggregate molecule.

Sanderson's geometric mean equalization principle for electronegativity has been used to approximate the reactivity parameter electronegativity of an aggregate from the corresponding properties of the constituent fragments.<sup>6–8</sup> The Sanderson scheme was later extended to determine the molecular hardness.<sup>9</sup> The other related methods for calculating aggregate hardness from fragment values are that of the average softness method introduced by Ghosh et al.<sup>10</sup> and the geometric mean principle of hardness introduced by Datta<sup>11</sup> and later extended by Chattaraj et al.<sup>11</sup>

Recently, an additivity scheme for the chemical potential and hardness has been introduced,<sup>12</sup> and the novelty of that scheme is that it involves the Fukui function of component fragments for the estimation of the global properties of the aggregates. In the present work, we have obtained workable expressions for the calculation of the hardness parameter from the relevant theoretical expression and have compared the reactivity parameters viz. chemical potential and hardness so obtained with the reference ab initio results and computed from the other additivity schemes currently available.<sup>10,11</sup>

## 2. Theoretical Background

As mentioned, within the density functional theory (DFT) formalism, the chemical potential and the hardness for an  $N$ -particle system with total energy  $E$  and external potential  $v(\mathbf{r})$  are defined as

$$\mu = \left( \frac{\partial E}{\partial N} \right)_{v(\mathbf{r})} = -\chi \quad (1)$$

and

$$\eta = \frac{1}{2} \left( \frac{\partial^2 E}{\partial N^2} \right)_{v(\mathbf{r})} = \frac{1}{2} \left( \frac{\partial \mu}{\partial N} \right)_{v(\mathbf{r})} \quad (2)$$

Using Koopman's theorem<sup>2</sup> and first-order finite difference approximation of eqs 1 and 2, we have

$$\mu \approx \frac{1}{2}(I + A) \approx \frac{1}{2}(\epsilon_L + \epsilon_H) \quad (3)$$

and

$$\eta \approx \frac{1}{2}(I - A) \approx \frac{1}{2}(\epsilon_L - \epsilon_H) \quad (4)$$

where  $I$  is the ionization potential and  $A$  is the electron affinity and  $\epsilon_H$  and  $\epsilon_L$  are the energies of the highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO), respectively.

**Sanderson's Additivity Scheme for  $\mu$  and  $\eta$ .** Sanderson's equalization principle states that the electronegativity of a molecule is given by the geometric mean of the electronegativities of the constituent isolated atoms or fragments. Accordingly, Sanderson's expression for the aggregate chemical potential is expressed as

$$\mu_n^s = - \left( \prod_i^n |\mu_i| \right)^{1/n} \quad (5)$$

where  $\mu_i$  is the chemical potential of the  $i$ th fragment. Differentiating Sanderson's expression for the aggregate  $\mu_n^s$  with respect to  $N$ ,<sup>9,12</sup> we obtain an expression for chemical hardness henceforth referred to as the Sanderson's relation for aggregate hardness:

$$\eta_n^s = \frac{\mu_n^s}{n} \sum_i \frac{\eta_i}{\mu_i} \quad (6)$$

where  $\eta_i$  is the hardness of the  $i$ th fragment. Note that the numerical factor  $1/2$  has been dropped from the original definitions.

**Ghosh Additivity Scheme for  $\eta$ .** The combination scheme proposed by Ghosh et al.<sup>10</sup> expresses the aggregate hardness as

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the reciprocal of the aggregate softness computed as the average of the softness of the constituent fragments. Accordingly, it is calculated as

$$\frac{1}{\eta_n^s} = \frac{1}{n} \sum_i^n \frac{1}{\eta_i} \quad (7)$$

**Datta Additivity Scheme for  $\eta$ .** Another aggregation scheme has been introduced by Datta<sup>11</sup> which expresses the hardness of the aggregate molecule as the geometric mean of the respective hardness of the constituent fragments as follows:

$$\eta_n^d = \left( \prod_i^n \eta_i \right)^{1/n} \quad (8)$$

**New Additivity Scheme for  $\mu$  and  $\eta$ .** A new additivity scheme for  $\mu$  and  $\eta$  was recently proposed,<sup>12</sup> it comes out from the Sanderson's expression for the chemical potential (see eq 5). The chemical potential is a function of  $N$  and a functional of the external potential  $v(\mathbf{r})$ , and differentiation of  $\mu_n^s$  with respect to  $N$  led to eq 6, whereas differentiation with respect to  $v(\mathbf{r})$  leads to

$$F_n^{\circ} = \int \left( \frac{\delta \mu_n^s}{\delta v} \right) d\mathbf{r} \approx \frac{\mu_n^s}{n} \sum_i^n \left( \frac{f_i}{\mu_i} \right) \quad (9)$$

where  $F_n^{\circ}$  is a dimensionless quantity that contains the fragment reactivities through the condensed Fukui function  $f_i$  associated with the  $i$ th constituent fragment of the molecule.<sup>12</sup> The Fukui function is a local property that has been defined as<sup>2</sup>

$$f(\mathbf{r}) = \left( \frac{\delta \rho(\mathbf{r})}{\delta N} \right)_v \quad (10)$$

which clearly integrates to unity. For removal or addition of charge from or to a neutral atom or molecule to obtain another ground state, it has been proved that<sup>13</sup>

$$f^-(\mathbf{r}) \approx [\rho_N(\mathbf{r}) - \rho_{N-1}(\mathbf{r})] \quad (11)$$

measures the reactivity toward a electrophilic attack, in this case it has been assumed that one electron was removed by the electrophile from our reference  $N$  electron system. On the other hand

$$f^+(\mathbf{r}) \approx [\rho_{N+1}(\mathbf{r}) - \rho_N(\mathbf{r})] \quad (12)$$

measures the reactivity toward a nucleophilic attack, here the assumption is that one electron is transferred from the nucleophile to the reference  $N$  electron system. Using the local density approximation (LDA)<sup>14</sup> for assuming uniform removal or addition of charge, the Fukui function can be approximated as

$$f^-(\mathbf{r}) \approx f^+(\mathbf{r}) \approx f(\mathbf{r}) \approx \frac{\rho(\mathbf{r})}{N} \quad (13)$$

After integration over, say, the  $i$ th fragment volume, we obtain a condensed Fukui function as

$$f_i = \frac{N_i}{N} \quad (14)$$

with  $\sum_i N_i = N$  as the total number of electrons in the aggregate and  $N_i$  is the number of electrons in the  $i$ th fragment. In the

context of the formation of a molecular aggregate from the constituent fragments, the Fukui function as defined in eq 14 will play the role of a weighting factor instead of a reactivity parameter.

Because  $F_n^{\circ}$  and  $f_i$  are dimensionless quantities, the dimensional analysis of the right-hand side of eq 9 shows that the sum containing the  $f_i/\mu_i$  terms should be of dimension (1/chemical potential). Therefore, we propose the chemical potential for a composite system be defined as

$$\frac{1}{\mu_n^t} = \sum_i^n \left( \frac{f_i}{\mu_i} \right) \quad (15)$$

The corresponding hardness is obtained by differentiating eq 15 with respect to  $N$  at constant external potential, which leads to

$$\eta_n^t = \left( \frac{\partial \mu_n^t}{\partial N} \right)_{v(\mathbf{r})} = (\mu_n^t)^2 \sum_i^n \frac{1}{(\mu_i)^2} [f_i \eta_i - h_i \mu_i] \quad (16)$$

where

$$h_i = \left( \frac{\partial f_i}{\partial N} \right)_{v(\mathbf{r})} \quad (17)$$

$h_i$  is a measure of the change in the fragment Fukui function because of the change in the total number of electrons. It is not possible to obtain an explicit evaluation of this reactivity descriptor. So, in the present work, for two component aggregates, we obtain an approximation by rewriting eq 17 as follows:

$$\begin{aligned} h_1 &\approx \frac{\Delta f_1}{\Delta N} \\ &\approx \frac{f_1(\text{product}) - f_1(\text{reactant})}{N(\text{product}) - N(\text{reactant})} \\ &\approx \frac{N_1/N(\text{product}) - 1}{N(\text{product}) - N_1} \\ &\approx -\frac{1}{N} \end{aligned} \quad (18)$$

where  $N(\text{product}) = N_1 + N_2$ , so that for bicomponent aggregates, the expression for hardness as given by eq 16 can be approximated as

$$\eta_2^t = \left( \frac{\partial \mu_2^t}{\partial N} \right)_{v(\mathbf{r})} \approx \frac{(\mu_2^t)^2}{N} \sum_{i=1}^2 \frac{1}{(\mu_i)^2} [N_i \eta_i + \mu_i] \quad (19)$$

We have also approximated eq 16 assuming  $f_i$  to remain constant as a fragment aggregates so that  $h_i = 0$ ; accordingly, we have

$$\eta_2^{t_0} = \left( \frac{\partial \mu_2^t}{\partial N} \right)_{v(\mathbf{r})} \approx \frac{(\mu_2^t)^2}{N} \sum_{i=1}^2 \frac{1}{(\mu_i)^2} [N_i \eta_i] \quad (20)$$

### 3. Computational Details

All calculations were performed at the RHF level of theory with the standard 6-311G\*\* basis set using the Gaussian 98 package.<sup>15</sup> Radical fragments calculations were carried out applying the UHF theory. The electronic chemical potential and

**TABLE 1: Electronic Chemical Potential at the HF/6-311G\*\* Level of Calculation for 46 Molecular Aggregates ( $\mu^\circ$ ) and the Corresponding Approximate Values Determined through the Use of Various Combinations Schemes<sup>a</sup>**

| no. | fragments  | aggregate  | $\mu^\circ$ | $\mu_2^s$ | $\mu_2^t$ |
|-----|--|--|-------------|-----------|-----------|
| 1   | CHO + CHO  | CHOCHO   | -0.1852     | -0.1324   | -0.1324   |
| 2   | CHO + OCl  | CHOCl  | -0.1891     | -0.1567   | -0.1613   |
| 3   | CHO + OH   | HCOOH  | -0.1554     | -0.1588   | -0.1495   |
| 4   | CHO + SH   | HCOSH  | -0.1391     | -0.1589   | -0.1581   |
| 5   | CN + CHO   | CNCHO  | -0.2223     | -0.1896   | -0.1737   |
| 6   | H + CHO  | HCHO   | -0.1493     | -0.0999   | -0.1264   |
| 7   | H + HS   | H <sub>2</sub> S                                   | -0.1267     | -0.1120   | -0.1758   |
| 8   | H + CN   | HCN  | -0.1677     | -0.1431   | -0.2290   |
| 9   | H + OCl  | HOCl   | -0.1733     | -0.1183   | -0.1757   |
| 10  | HS + OH  | HSOH   | -0.1317     | -0.1907   | -0.1907   |
| 11  | HS + NO  | HSNO   | -0.1571     | -0.1697   | 0.1697    |
| 12  | SN + OH  | SNOH   | -0.1746     | -0.1846   | -0.1819   |
| 13  | CH <sub>3</sub> + CHO  | CH <sub>3</sub> CHO                                | -0.1326     | -0.1283   | -0.1292   |
| 14  | CH <sub>3</sub> + CHS  | CH <sub>3</sub> CHS                                | -0.1397     | -0.1362   | -0.1413   |
| 15  | CH <sub>3</sub> + HS   | CH <sub>3</sub> SH                                 | -0.1125     | -0.1540   | -0.1610   |
| 16  | CH <sub>3</sub> + CH <sub>3</sub>  | CH <sub>3</sub> CH <sub>3</sub>                    | -0.1644     | -0.1243   | -0.1243   |
| 17  | CH <sub>3</sub> + CN   | CH <sub>3</sub> CN                                 | -0.1662     | -0.1837   | -0.1829   |
| 18  | CH <sub>3</sub> + OH   | CH <sub>3</sub> OH                                 | -0.1447     | -0.1539   | -0.1505   |
| 19  | CH <sub>2</sub> + CH <sub>2</sub>  | C <sub>2</sub> H <sub>4</sub>                      | -0.1064     | -0.1657   | -0.1657   |
| 20  | CH <sub>3</sub> + C <sub>2</sub> H <sub>5</sub>                                  | C <sub>3</sub> H <sub>8</sub>                      | -0.1560     | -0.1124   | -0.1085   |
| 21  | CH <sub>2</sub> + C <sub>2</sub> H <sub>6</sub>                                  | C <sub>3</sub> H <sub>8</sub>                      | -0.1560     | -0.1651   | -0.1648   |
| 22  | C <sub>2</sub> H <sub>5</sub> + CHO  | C <sub>2</sub> H <sub>5</sub> CHO                  | -0.1379     | -0.1160   | -0.1141   |
| 23  | CH <sub>2</sub> + C <sub>2</sub> H <sub>4</sub> O                                | C <sub>2</sub> H <sub>5</sub> CHO                  | -0.1379     | -0.1482   | -0.1396   |
| 24  | C <sub>2</sub> H <sub>5</sub> + OH   | C <sub>2</sub> H <sub>5</sub> OH                   | -0.1414     | -0.1392   | -0.1213   |
| 25  | CH <sub>2</sub> + CH <sub>3</sub> OH   | C <sub>2</sub> H <sub>5</sub> OH                   | -0.1414     | -0.1549   | -0.1506   |
| 26  | C <sub>2</sub> H <sub>5</sub> + CHS  | C <sub>2</sub> H <sub>5</sub> CHS                  | -0.1367     | -0.1232   | -0.1245   |
| 27  | CH <sub>2</sub> + CH <sub>3</sub> CHS  | C <sub>2</sub> H <sub>5</sub> CHS                  | -0.1367     | -0.1522   | -0.1442   |
| 28  | C <sub>2</sub> H <sub>5</sub> + CN   | C <sub>2</sub> H <sub>5</sub> CN                   | -0.1652     | -0.1662   | -0.1395   |
| 29  | CH <sub>3</sub> + CH <sub>2</sub> CN   | C <sub>2</sub> H <sub>5</sub> CN                   | -0.1652     | -0.1660   | -0.1610   |
| 30  | C <sub>3</sub> H <sub>7</sub> + CHO  | C <sub>3</sub> H <sub>7</sub> CHO                  | -0.1318     | -0.1203   | -0.1170   |
| 31  | CH <sub>3</sub> CN + C <sub>2</sub> H <sub>4</sub>                               | C <sub>3</sub> H <sub>7</sub> CN                   | -0.1612     | 0.1357    | -0.1342   |
| 32  | CH <sub>2</sub> CN + C <sub>2</sub> H <sub>5</sub>                               | C <sub>3</sub> H <sub>7</sub> CN                   | -0.1612     | -0.1330   | -0.1344   |
| 33  | C <sub>2</sub> H <sub>5</sub> + C <sub>2</sub> H <sub>5</sub>                    | C <sub>4</sub> H <sub>10</sub>                     | -0.1490     | -0.1017   | -0.1017   |
| 34  | C <sub>2</sub> H <sub>4</sub> + C <sub>2</sub> H <sub>6</sub>                    | C <sub>4</sub> H <sub>10</sub>                     | -0.1490     | -0.1323   | -0.1309   |
| 35  | CH <sub>3</sub> CHS + CH <sub>3</sub> CHCH <sub>2</sub>                          | C <sub>4</sub> H <sub>9</sub> CHS                  | -0.1372     | -0.1174   | -0.1185   |
| 36  | C <sub>2</sub> H <sub>5</sub> CHS + C <sub>2</sub> H <sub>4</sub>                | C <sub>4</sub> H <sub>9</sub> CHS                  | -0.1372     | -0.1206   | -0.1264   |
| 37  | C <sub>2</sub> H <sub>4</sub> + C <sub>2</sub> H <sub>4</sub>                    | CH <sub>3</sub> CH = CHCH <sub>3</sub>             | -0.0941     | -0.1064   | -0.1064   |
| 38  | CH <sub>2</sub> CH <sub>3</sub> + CHCH <sub>3</sub>                              | CH <sub>3</sub> CH = CHCH <sub>3</sub>             | -0.0941     | -0.1278   | -0.1097   |
| 39  | CH <sub>2</sub> + C <sub>2</sub> H <sub>4</sub>                                  | CH <sub>3</sub> CHCH <sub>2</sub>                  | -0.0986     | -0.1328   | -0.1208   |
| 40  | CH <sub>2</sub> + C <sub>3</sub> H <sub>7</sub> CHO                              | C <sub>2</sub> H <sub>5</sub> CHCHOCH <sub>3</sub> | -0.1348     | -0.1478   | -0.1365   |
| 41  | C <sub>2</sub> H <sub>4</sub> + C <sub>2</sub> H <sub>5</sub> CHO                | C <sub>2</sub> H <sub>5</sub> CHCHOCH <sub>3</sub> | -0.1348     | -0.1211   | -0.1255   |
| 42  | CH <sub>3</sub> CHCH <sub>2</sub> + CH <sub>3</sub> CHCH <sub>2</sub>            | C <sub>6</sub> H <sub>12</sub>                     | -0.0850     | -0.0986   | -0.0986   |
| 43  | C <sub>2</sub> H <sub>4</sub> + CH <sub>2</sub> C(CH <sub>3</sub> ) <sub>2</sub> | C <sub>6</sub> H <sub>12</sub>                     | -0.0850     | -0.1019   | -0.1004   |
| 44  | CH <sub>2</sub> + C <sub>3</sub> H <sub>8</sub>                                  | CH(CH <sub>3</sub> ) <sub>3</sub>                  | -0.1529     | -0.1608   | -0.1582   |
| 45  | CH <sub>3</sub> + C <sub>3</sub> H <sub>7</sub>                                  | CH(CH <sub>3</sub> ) <sub>3</sub>                  | -0.1529     | -0.1166   | -0.1130   |
| 46  | CH <sub>2</sub> + CH <sub>2</sub> CHCHO  | CH <sub>3</sub> CHCHCHO                            | -0.1399     | -0.1601   | -0.1569   |
|     | $\sigma^b$   |  |             | 0.0278    | 0.0283    |

<sup>a</sup> All values are in au. <sup>b</sup>  $\sigma$  is the root-mean-squared deviation from ab initio values.

the molecular hardness were computed using eqs 3 and 4, respectively. We have studied a total of 46 unique combinations to yield 34 aggregates composed from pairs of 37 fragments in all.

The chemical potentials of the aggregates were determined using Sanderson's and our own expressions that are given in eqs 5 and 15, respectively. For a comparative study of the relative merits of the various estimation schemes for the hardness parameter, calculations have been carried out using eq 6 (Sanderson) and methods of arithmetic average of softness<sup>10</sup> (eq 7) and the geometric mean principle of hardness<sup>11</sup> (eq 8). These results are compared to those determined using eqs 19 and 20.

#### 4. Results and Discussion

The estimated chemical potential and hardness values in atomic units are compared with the corresponding ab initio results for the aggregates denoted as  $\mu^\circ$  and  $\eta^\circ$  given in Tables 1 and 2, respectively. The root-mean-squared deviations of

chemical potentials and hardness parameters calculated as per the schemes discussed here from the corresponding ab initio values for the aggregates are also quoted in Tables 1 and 2. Evidently, for the systems studied in this work, the new combination scheme for the hardness introduced here as well as the chemical potential yields results very much comparable to those obtained from other relevant schemes.

The results quoted in Table 1 show that in most of the cases the chemical potential  $\mu_2^t$  as introduced by Gutiérrez-Oliva et al.<sup>12</sup> yields values closer to the ab initio results as compared to the estimates obtained from eq 5 as proposed by Sanderson. The main reason for this marked improvement may be due to the introduction of  $f_i$  in the calculation of aggregate property effectively reflecting the local electronic population within that aggregate. We note that for lighter aggregates composed by a few atoms the Sanderson's expression for chemical potential  $\mu_2^s$  yields closer agreement with the ab initio results. This may be due to the fact that for smaller constituents with fewer atoms the local properties and charge rearrangements play a relatively

**TABLE 2: Molecular Hardness at the HF/6-311G\*\* Level of Calculation for 46 Molecular Aggregates ( $\eta^\circ$ ) and the Corresponding Approximate Values Determined through the Use of Various Combinations Schemes<sup>a</sup>**

| no.        | $\eta^\circ$ | $\eta_2^s$ | $\eta_2^g$ | $\eta_2^d$ | $\eta_2^t$ | $\eta_2^0$ |
|------------|--------------|------------|------------|------------|------------|------------|
| 1          | 0.2396       | 0.2453     | 0.2453     | 0.2453     | 0.2365     | 0.2453     |
| 2          | 0.2737       | 0.2563     | 0.2538     | 0.2540     | 0.2523     | 0.2607     |
| 3          | 0.3138       | 0.2785     | 0.2759     | 0.2781     | 0.2563     | 0.2683     |
| 4          | 0.2608       | 0.2249     | 0.2119     | 0.2138     | 0.2220     | 0.2317     |
| 5          | 0.2793       | 0.2636     | 0.2486     | 0.2487     | 0.2620     | 0.2741     |
| 6          | 0.2866       | 0.3736     | 0.3109     | 0.3226     | 0.2634     | 0.2841     |
| 7          | 0.2582       | 0.3960     | 0.2591     | 0.2813     | 0.2458     | 0.2776     |
| 8          | 0.3160       | 0.4689     | 0.3163     | 0.3271     | 0.3822     | 0.4456     |
| 9          | 0.2700       | 0.4166     | 0.3248     | 0.3341     | 0.2931     | 0.3152     |
| 10         | 0.2431       | 0.2509     | 0.2343     | 0.2424     | 0.2164     | 0.2311     |
| 11         | 0.2465       | 0.2276     | 0.2162     | 0.2190     | 0.2204     | 0.2311     |
| 12         | 0.2019       | 0.2474     | 0.2320     | 0.2405     | 0.2062     | 0.2174     |
| 13         | 0.2878       | 0.2530     | 0.2524     | 0.2525     | 0.2406     | 0.2515     |
| 14         | 0.2023       | 0.2270     | 0.2163     | 0.2195     | 0.2046     | 0.2138     |
| 15         | 0.2448       | 0.2363     | 0.2172     | 0.2202     | 0.2245     | 0.2377     |
| 16         | 0.3236       | 0.2600     | 0.2600     | 0.2600     | 0.2462     | 0.2600     |
| 17         | 0.3024       | 0.2774     | 0.2560     | 0.2560     | 0.2800     | 0.2979     |
| 18         | 0.3002       | 0.2883     | 0.2850     | 0.2863     | 0.2720     | 0.2888     |
| 19         | 0.2725       | 0.2243     | 0.2243     | 0.2243     | 0.2036     | 0.2243     |
| 20         | 0.3131       | 0.2562     | 0.2553     | 0.2553     | 0.2472     | 0.2553     |
| 21         | 0.3131       | 0.2742     | 0.2650     | 0.2694     | 0.2808     | 0.2934     |
| 22         | 0.2757       | 0.2505     | 0.2480     | 0.2480     | 0.2460     | 0.2531     |
| 23         | 0.2757       | 0.2612     | 0.2521     | 0.2541     | 0.2707     | 0.2790     |
| 24         | 0.2961       | 0.2868     | 0.2793     | 0.2812     | 0.2688     | 0.2774     |
| 25         | 0.2961       | 0.2654     | 0.2568     | 0.2595     | 0.2708     | 0.2820     |
| 26         | 0.2031       | 0.2283     | 0.2131     | 0.2155     | 0.2274     | 0.2339     |
| 27         | 0.2031       | 0.2132     | 0.2128     | 0.2130     | 0.1997     | 0.2065     |
| 28         | 0.2980       | 0.2820     | 0.2514     | 0.2514     | 0.2874     | 0.2962     |
| 29         | 0.2980       | 0.2633     | 0.2576     | 0.2605     | 0.2704     | 0.2815     |
| 30         | 0.2820       | 0.2448     | 0.2438     | 0.2438     | 0.2394     | 0.2452     |
| 31         | 0.2959       | 0.2575     | 0.2457     | 0.2458     | 0.2611     | 0.2684     |
| 32         | 0.2959       | 0.2913     | 0.2867     | 0.2871     | 0.2903     | 0.2976     |
| 33         | 0.3078       | 0.2508     | 0.2508     | 0.2508     | 0.2448     | 0.2508     |
| 34         | 0.3078       | 0.2996     | 0.2959     | 0.2970     | 0.2950     | 0.3025     |
| 35         | 0.2008       | 0.2411     | 0.2285     | 0.2304     | 0.2414     | 0.2457     |
| 36         | 0.2008       | 0.2440     | 0.2328     | 0.2353     | 0.2292     | 0.2340     |
| 37         | 0.2504       | 0.2725     | 0.2725     | 0.2725     | 0.2660     | 0.2725     |
| 38         | 0.2531       | 0.2566     | 0.2419     | 0.2426     | 0.2621     | 0.2682     |
| 39         | 0.2624       | 0.2600     | 0.2461     | 0.2473     | 0.2646     | 0.2740     |
| 40         | 0.2732       | 0.2581     | 0.2499     | 0.2515     | 0.2719     | 0.2772     |
| 41         | 0.2732       | 0.2762     | 0.2741     | 0.2741     | 0.2732     | 0.2787     |
| 42         | 0.2370       | 0.2624     | 0.2624     | 0.2624     | 0.2583     | 0.2624     |
| 43         | 0.2370       | 0.2910     | 0.2883     | 0.2888     | 0.2930     | 0.2988     |
| 44         | 0.3060       | 0.2701     | 0.2614     | 0.2650     | 0.2851     | 0.2943     |
| 45         | 0.3060       | 0.2511     | 0.2506     | 0.2510     | 0.2405     | 0.2469     |
| 46         | 0.2430       | 0.2361     | 0.2350     | 0.2353     | 0.2346     | 0.2427     |
| $\sigma^b$ |              | 0.0491     | 0.0338     | 0.0340     | 0.0344     | 0.0340     |

<sup>a</sup> All values are in au. <sup>b</sup>  $\sigma$  is the root-mean-squared deviation from *ab initio* values.

minor role, and in such cases the LDA fukui function used in  $\mu_2^t$  may overestimate the local effects. In fact, this view is borne out further by the observation that for heavier aggregates composed of more atoms  $\mu_2^t$  consistently yields better agreement with the *ab initio* results for the aggregates.

In Table 2 are quoted the values of molecular hardness estimated following different schemes along with the corresponding *ab initio* results. It can be noticed that the expression  $\eta_2^t$  gives consistently better results than  $\eta_2^s$  and  $\eta_2^d$  and compares favorably with  $\eta_2^g$ . As before, the incorporation of the local fukui function appears to be the primary reason for this significant improvement in the estimated values. It should be mentioned that in general the estimates of the chemical potentials are better compared to the hardness estimates, with the fundamental reason for this being the second order differential involved in the expression for hardness.

Although overall numerical results obtained using eqs 15 and 16 are quite comparable to those determined through other additivity schemes, the merit of  $\mu_n^t$  and  $\eta_n^t$  is basically that they involve local quantities that act as weighting factors in the combination of the fragments to form the molecule.

## Concluding Remarks

In the present work, we have established that the proposed chemical potential additivity scheme incorporating the fragment condensed fukui function yields consistently better results for the estimation of aggregate chemical potential compared to those obtained following Sanderson's scheme. In fact, this brings out the vital importance of the *weight* of the constituent fragments within the composite system. Also, two new related schemes for the estimation of hardness have been introduced which are comparable with the results from Ghosh's scheme using the inverse of softness average. Most importantly, the explicit albeit approximate evaluation of the fragment hardness response function to the change in the external potential during bifragment aggregation has enabled the incorporation of the local fukui function in the expression for hardness thus adding fragments' weighing factors not considered in other schemes. The consequent encouraging results open up the possibility that better approximations of the fukui function can enhance the quality of prediction of global aggregate properties from the constituent fragment properties.

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